Software Package to Calculate the Effects of Core hole and Surface Excitations on XPS and AES Sven Tougaard , Francisco Yubero

QUEELS-XPS software package Semi-classical dielectric response model

- 1. Intro: REELS Semi-classical dielectric response model
- 2. The XPS- Semiclassical dielectric response model Intrinsic and extrinsic excitations
- 3. The QUEELS-XPS software package
- 4. The validity of the model quantitative comparison with experiments

# Intro: REELS Semi-classical dielectric response model QUEELS-ε(k,ω)-REELS software package

2. The XPS- Semiclassical dielectric response model Extrinsic and intrinsic excitations

- 3. The QUEELS-XPS software Package
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#### Dielectric response model for REELS

Tougaard and ChorkendorffPhys Rev B35, 6570 (1987)Yubero and TougaardPhys Rev B46, 2486 (1992)Yubero, Sanz, Ramskov, Tougaard Phys Rev B53, 9719 (1996)





Low energies are necessary

Surface effects and k-dispersion effects are important

Quantitative determination of dielectric function optical properties of nano-meter thin films

#### Semiclassical dielectric response model for REELS



Effective Peross section J. M. SANZ, B. RAMSKOV, AND S. TOUGAARD

Each term can be found from the following expressions:

# $\epsilon(\mathbf{k},\omega)$ is the only unknown $\left[ \sum_{k_{\perp}-\Omega_{i}=1}^{\epsilon(k_{\perp}-\Omega_{i})-i} \frac{e^{-i(k_{\perp}-\Omega_{i})a}}{k_{\perp}-\Omega_{i}} \right]$

$$\mathcal{K}_{\text{eff}}^{m_i} = \operatorname{Re}\left\{\frac{1}{(2\pi)^4 \hbar^2 \omega x} \int d\mathbf{k} \left(k_{\perp} + k_{\parallel} \frac{\omega}{v_{\perp i}}\right) e^{ik_{\perp}a} \Phi_{\text{ind}}^{m}(\mathbf{k}, \omega) \left[\frac{1 - e^{-i(k_{\perp} - \Omega_i)a}}{k_{\perp} - \Omega_i}\right]\right\},\tag{18}$$

$$K_{\text{eff}}^{V_o} = \operatorname{Re}\left(\frac{-2ie}{(2\pi)^4 \hbar^2 \omega x} \int d\mathbf{k} \left(k_{\perp} - k_{\parallel} \frac{v_{\parallel o}}{v_{\perp o}}\right) e^{ik_{\perp}a} \Phi_{\text{ind}}^V(\mathbf{k}, \omega) \left[\pi \delta(k_{\perp} + \Omega_o) - i \frac{e^{-i(k_{\perp} + \Omega_o)a}}{k_{\perp} + \Omega_o}\right]\right),\tag{19}$$

$$\frac{M_o}{\text{eff}} = \text{Re} \left\{ \frac{2e}{(2\pi)^4 \hbar^2 \omega x} \int d\mathbf{k} \left( k_\perp - k_\parallel \frac{v_{\parallel o}}{v_{\perp o}} \right) e^{i k_\perp a} \Phi^M_{\text{ind}}(\mathbf{k}, \omega) \left[ \frac{1 - e^{-i (k_\perp + \Omega_o)}}{k_\perp + \Omega_o} \right] \right\}.$$
(20)

Equations (16)–(20) give the solution to the problem of finding  $K_{\text{eff}}$  for a general REELS geometry. However, in the following we will consider approximations that allow some integrals to be done analytically.

The integrals in momentum are given by the conservation laws for energy and momentum. Unfortunately a complete analytical integration is not possible. However, with cylindrical coordinates  $(dk=2\pi dk_{\parallel}dk_{\perp})$  and extending the limits of integration over  $k_{\perp}$  to  $-\infty < k_{\perp} < +\infty$ , the integral over  $k_{\perp}$  can be done analytically. With this approximation we do not expect large errors because the main contribution to the integral comes from small values of k, and the integrand decreases roughly as  $1/k^9$  for  $k \to \pm\infty$  using the functional shape for the ELF as in Eq. (30) (see below). Assuming again  $\epsilon = \epsilon(k_{\parallel}, \omega)$ , it is found that

$$K_{\text{eff}}^{V_i} = \text{Re}\left\{\frac{e}{\hbar^2 \omega \pi x} \frac{v_{\perp i} + i v_{\parallel i}}{v_{\perp i}} \int dk_{\parallel} \frac{\epsilon - 1}{\epsilon + 1} \frac{k_{\parallel}(k_{\parallel} + i\Omega_i)}{k_{\parallel}^2 + \Omega_i^2} e^{i\Omega_i a} [F^{MM}(k_{\parallel}) + F^{MV}(k_{\parallel})]\right\},\tag{21}$$

$$K_{\text{eff}}^{V_{o}} = \operatorname{Re}\left\{\frac{-e}{\hbar^{2}\omega\pi x}\frac{\upsilon_{\perp o} - i\upsilon_{\parallel o}}{\upsilon_{\perp o}}\int dk_{\parallel}\frac{\epsilon-1}{\epsilon+1}\frac{k_{\parallel}(k_{\parallel}-i\Omega_{o})}{k_{\parallel}^{2}+\Omega_{o}^{2}}e^{-i\Omega_{o}a}[F^{MM}(k_{\parallel})+F^{MV}(k_{\parallel})]\right\},\tag{22}$$

$$\begin{split} \mathcal{K}_{\text{eff}}^{M_{i}} = & \operatorname{Re} \left\{ \frac{-2ie^{2}}{\hbar^{2} \pi v_{\perp i}^{2} x} \int dk_{\parallel} \left( \frac{1}{\epsilon} - 1 \right) \frac{k_{\parallel}}{k_{\parallel}^{2} + \Omega_{i}^{2}} \right\} + \operatorname{Re} \left\{ \frac{-2e^{2}}{\hbar^{2} \pi v_{\perp o} \omega_{w}} \int dk_{\parallel} \left( \frac{1}{\epsilon} - 1 \right) \frac{e^{i(\Omega_{i} + \Omega_{o})} - 1}{\Omega_{i} + \Omega_{o}} \frac{k_{\parallel}}{k_{\parallel}^{2} + \Omega_{o}^{2}} \left[ k_{\parallel} \left( \frac{v_{\perp i}}{v_{\perp i}} + \frac{v_{\parallel o}}{v_{\perp o}} \right) \right] \\ & - \frac{\omega}{v_{\perp o}} \right] \right\} + \operatorname{Re} \left\{ \frac{e}{\hbar^{2} \pi \omega x} \frac{v_{\perp i} - iv_{\parallel i}}{v_{\perp i}} \int dk_{\parallel} \left( \frac{1}{\epsilon} - 1 \right) \frac{k_{\parallel}(k_{\parallel} - i\Omega_{i})}{k_{\parallel}^{2} + \Omega_{i}^{2}} \left( e^{-k_{l}a} - e^{i\Omega_{i}a} \right) \rho_{1}^{M}(k_{\perp} = ik_{\parallel}) \right\} \\ & + \operatorname{Re} \left\{ \frac{e}{\hbar^{2} \pi \omega x} \frac{v_{\perp i} - iv_{\parallel i}}{v_{\perp i}} \int dk_{\parallel} \left( \frac{1}{\epsilon} - 1 \right) \frac{k_{\parallel}(k_{\parallel} + i\Omega_{i})}{k_{\parallel}^{2} + \Omega_{i}^{2}} \left( e^{k_{l}a} - e^{i\Omega_{i}a} \right) \rho_{2}^{M}(k_{\perp} = -ik_{\parallel}) \right\} \\ & + \operatorname{Re} \left\{ \frac{e}{\hbar^{2} \pi \omega x} \frac{v_{\perp i} - iv_{\parallel i}}{2v_{\perp i}} \int dk_{\parallel} \left( \frac{(\epsilon + 2)(\epsilon - 1)}{\epsilon(\epsilon + 1)} \frac{F^{MM}(k_{\parallel}) - \frac{\epsilon - 1}{\epsilon + 1}}{F^{MV}(k_{\parallel})} \right) \frac{k_{\parallel}(k_{\parallel} - i\Omega_{i})}{k_{\parallel}^{2} + \Omega_{i}^{2}} \left( e^{-k_{\parallel}a} - e^{i\Omega_{i}a} \right) \rho_{2}^{M}(k_{\perp} = -ik_{\parallel}) \right\} \\ & + \operatorname{Re} \left\{ \frac{e}{\hbar^{2} \pi \omega x} \frac{v_{\perp i} - iv_{\parallel i}}{2v_{\perp i}} \int dk_{\parallel} \left( \frac{(\epsilon + 2)(\epsilon - 1)}{\epsilon(\epsilon + 1)} \frac{F^{MM}(k_{\parallel}) - \frac{\epsilon - 1}{\epsilon + 1}} \frac{F^{MV}(k_{\parallel})}{k_{\parallel}^{2} + \Omega_{i}^{2}} \left( e^{-k_{\parallel}a} - e^{i\Omega_{i}a} \right) \right\}, \quad (23) \\ & \mathcal{K}_{\text{eff}}^{M_{0}} = \operatorname{Re} \left\{ \frac{-2ie^{2}}{\hbar^{2} \pi w_{\perp}^{2} \alpha_{w}} \int dk_{\parallel} \left( \frac{1}{\epsilon} - 1 \right) \frac{k_{\parallel}}{k_{\parallel}^{2} + \Omega_{o}^{2}} \right\} + \operatorname{Re} \left\{ \frac{2e^{2}}{\hbar^{2} \pi w_{\perp \perp \omega \omega}} \int dk_{\parallel} \left( \frac{1}{\epsilon} - 1 \right) \frac{k_{\parallel}}{k_{\parallel}^{2} + \Omega_{o}^{2}}} \left( e^{-k_{\parallel}a} - e^{-i\Omega_{o}a} \right) \rho_{1}^{M}(k_{\perp} - ik_{\parallel}) \right\} \\ & - \frac{\omega}{w_{\perp \perp}}} \right\} + \operatorname{Re} \left\{ \frac{-e}{\hbar^{2} \pi \omega w} \frac{v_{\perp o} + iv_{\parallel o}}{v_{\perp o}} \int dk_{\parallel} \left( \frac{1}{\epsilon} - 1 \right) \frac{k_{\parallel}(k_{\parallel} + i\Omega_{o})}{k_{\parallel}^{2} + \Omega_{o}^{2}}} \left( e^{-k_{\parallel}a} - e^{-i\Omega_{o}a} \right) \rho_{1}^{M}(k_{\perp} - ik_{\parallel}) \right\} \\ & + \operatorname{Re} \left\{ \frac{-e}{\hbar^{2} \pi \omega w} \frac{v_{\perp o} + iv_{\parallel o}}{v_{\perp o}} \int dk_{\parallel} \left( \frac{1}{\epsilon} - 1 \right) \frac{k_{\parallel}(k_{\parallel} + i\Omega_{o})}{k_{\parallel}^{2} + \Omega_{o}^{2}}} \left( e^{-k_{\parallel}a} - e^{-i\Omega_{o}a} \right) \rho_{1}^{M}(k_{\parallel} - ik_{\parallel}) \right\} \\ & + \operatorname{Re} \left\{ \frac{-e}{\hbar^{2} \pi \omega w} \frac{v_{\perp o} + iv_{\parallel o}}{v_{\perp o}}$$

Notice that  $K_{\text{eff}}^{v_0,M_0}$  can be obtained from  $K_{\text{eff}}^{v_lM_l}$  by making the changes  $v_{\perp o} \rightarrow -v_{\perp i}$ ,  $v_{\perp j} \rightarrow -v_{\perp o}$ ,  $v_{\parallel i} \rightarrow v_{\parallel o}$ , and  $v_{\parallel o} \rightarrow v_{\parallel i}$ . Note also that in general  $K_{\text{eff}}(E_0,\hbar,\omega,a,\theta_0,\theta_j) \neq K_{\text{eff}}(E_0,\hbar,\omega,a,\theta_0,\theta_l)$ , i.e., the losses experienced by the electron depend not only on the line trajectory followed but also on the sense in which it is done. However, in practice the difference is very small. We must discuss now the limits in which the present model can be compared with model B in Ref. 9. That model is valid for normal incidence and exit angles of the electrons. This situation is given in the present model by considering  $v_{\parallel i} = v_{\parallel o} = 0$  and  $v_{\perp i} = v_{\perp o} = v$ , which gives 5

# Compare model calculation to REELS-experiment

Determination of the dielectric function  $\epsilon(k,\omega)$ 

QUEELS-  $\epsilon(k,\omega)$ -REELS software was developed to easily do this

(17)

Application to Si-gate oxides

Tahir et al. J. Appl. Phys 106 (2009) 084108



#### REELS exp comp to theory



Im [-1/ε(k,ω)]



Tahir et al. J. Appl. Phys 106 (2009) 084108

$$\varepsilon(\omega) = \varepsilon_1(\omega) + i\varepsilon_2(\omega)$$





QUEELS- ε(k,ω)-REELS software package

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#### **Photo excitation**

Intrinsic and extrinsic excitations:



Sudden creation of a core hole Intrinsic Shake-up excitations



Moving electron Extrinsic Energy loss

The contributions interfere and cannot be separated



Effective cross section

 $\mathbf{K}_{eff} = \mathbf{K}_{M1} + \mathbf{K}_{M2} + \mathbf{K}_{M3} + \mathbf{K}_{V} + \mathbf{K}_{M3} + \mathbf{K}_{V} + \mathbf{K}_{M1} + \frac{1}{2\pi E_{\perp} a_{0} x \omega_{\perp}} \operatorname{Re}\left[i \int k_{\parallel} dk_{\parallel} \frac{(1/\epsilon - 1)}{(\Omega_{0}^{2} + k_{\parallel}^{2})^{2}} \left\{\left(i + \frac{\nu_{\parallel}}{\nu_{\perp}}\right)(\Omega_{0} + ik_{\parallel})^{2}(e^{-i\Omega_{0}a - k_{\parallel}a} - 1) + \left(i - \frac{\nu_{\parallel}}{\nu_{\perp}}\right)(e^{-i\Omega_{0}a} - e^{-k_{\parallel}a})(e^{i\Omega_{0}a} - e^{-k_{\parallel}a})(\Omega_{0}^{2} + k_{\parallel}^{2})\right\}\right], \quad (13)$ 

$$K_{M2} = \frac{-1}{2\pi E_{\perp} a_0 x \omega_{\perp}} \operatorname{Re} \left[ \left( \pi \delta(\omega_{\perp}) + \frac{i}{\omega_{\perp}} \right) \int k_{\parallel} dk_{\parallel} \frac{(1/\epsilon - 1)}{\Omega_0^2 + k_{\parallel}^2} \left\{ \left( i + \frac{\nu_{\parallel}}{\nu_{\perp}} \right) (\Omega_0 + ik_{\parallel}) (1 - e^{-k_{\parallel} a - i\Omega_0 a}) + \left( -i + \frac{\nu_{\parallel}}{\nu_{\perp}} \right) (\Omega_0 - ik_{\parallel}) (e^{-2k_{\parallel} a} - e^{-k_{\parallel} a - i\Omega_0 a}) \right\} \right],$$
(14)

$$K_{M3} = \frac{-1}{2\pi E_{\perp}a_{0}x\omega_{\perp}} \operatorname{Re}\left[i\int k_{\parallel}dk_{\parallel} \frac{(-i+\nu_{\parallel}/\nu_{\perp})(\Omega_{0}-ik_{\parallel})(e^{-k_{\parallel}a}-e^{-i\Omega_{0}a})}{\Omega_{0}^{2}+k_{\parallel}^{2}} \left\{\frac{(\epsilon-1)(\epsilon+2)}{\epsilon(\epsilon+1)}\left[(e^{i\Omega_{0}a}-e^{-k_{\parallel}a})\left(\frac{\Omega_{0}+ik_{\parallel}}{\Omega_{0}^{2}+k_{\parallel}^{2}}\right) -i\left(\pi\delta(\omega_{\perp})+\frac{i}{\omega_{\perp}}\right)e^{-k_{\parallel}a}\right] + \frac{(\epsilon-1)}{(\epsilon+1)}\left(\frac{\Omega_{0}-ik_{\parallel}}{\Omega_{0}^{2}+k_{\parallel}^{2}}\right)e^{i\Omega_{0}a}\right\}\right],$$
(15)

$$K_{V} = \frac{1}{2\pi E_{\perp}a_{0}x\omega_{\perp}} \operatorname{Re}\left[\int k_{\parallel}dk_{\parallel} \frac{(1-\epsilon)}{(1+\epsilon)} \frac{(1-i\nu_{\parallel}/\nu_{\perp})(\Omega_{0}+ik_{\parallel})e^{-i\Omega_{0}a}}{\Omega_{0}^{2}+k_{\parallel}^{2}} \left\{ \frac{2ik_{\parallel}e^{i\Omega_{0}a}-(\Omega_{0}+ik_{\parallel})e^{-k_{\parallel}a}}{\Omega_{0}^{2}+k_{\parallel}^{2}} -i\left(\pi\delta(\omega_{\perp})+\frac{i}{\omega_{\perp}}\right)e^{-k_{\parallel}a}\right\}\right],$$

$$-i\left(\pi\delta(\omega_{\perp})+\frac{i}{\omega_{\perp}}\right)e^{-k_{\parallel}a}\right\}],$$

$$QUEELS-XPS$$
software package

 $\epsilon(k,\omega)$  is the only unknown

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#### QUEELS-XPS software

#### XPS-Simulation

File Tools Tools exclusively for the L\_K spectrum View Copyright

B+ XPS Simulation		
	Calculates XPS spectra	Information
Program Files Quases-Tougaard QUEELS-XPS Simulation Data X: .00, y: 0.00E+00	Exit Set ELF oscillators Calculate set of cross sections Calculate averaged cross section Keff,av for XPS	The software package calculates 1. Effective cross sections for photoelectrons excited at 65 different depths (eq.(2) in ref.[1]). 2. From the cross sections calculated in 1., the XPS - energy loss spectrum is calculated for a film of any thickness. The material is characterized by its dielectric response function ELF. The calculation take into account the effect of the core
	Geometry eq (X0)	The calculations build on the theory in 1. AC Simonsen, F Yubero, and S Tougaard Phys. Rev. B56, p. 1612 (1997) Examples of its practical use can be found in: 2. F Yubero and S Tougaard Phys Rev B71, 045414 (2005) 3. F Yubero, L Kover, W Drube, T Eickhoff, S Tougaard, Surf. Sci. 592 p. 1 (2005) 4. Z Berenyi et al, J Electr. Spectr and Rel. Phen. 135, p. 177 (2004) 5. L. Kover et al. Surf. Interf. Anal. 38, p. 569 (2006) 6. S Tougaard, F Yubero Surf Interf Anal. 36, p. 824 (2004) Copyright (c) 2001-2010 Sven Tougaard, University of Southern Denmark and Francisco Yubero ICMS, Seville, Spain The software is free for non commercial use 15







> XPS Simulation			
⊒ c: [Vista] ]C:\	ELF-SI EnergyLossFunction.dll	Calculates XPS spectra	Information
Program Files Cuases-Tougaard QUEELS-XPS Simulation Data		Exit	The software package calculates 1. Effective cross sections for photoelectrons excited at 65 different depths (eq.(2) in ref.[1]).
		Set ELF oscillators	<ol> <li>From the cross sections calculated in 1., the XPS - energy loss spectrum is calculated for a film of any thickness.</li> <li>The material is characterized by its dielectric</li> </ol>
(: .00, y: 0.00E+00		Calculate averaged cross section Keff,av for XPS	The calculation take into account the effect of the core
		Geometry e t q (X0)	The calculations build on the theory in 1. AC Simonsen, F Yubero, and S Tougaard Phys. Rev. B56, p. 1612 (1997) Examples of its practical use can be found in: 2. F Yubero and S Tougaard Phys Rev B71, 045414 (2005) 3. F Yubero, L Kover, W Drube, T Eickhoff, S Tougaard, Surf. Sci. 592 p. 1 (2005) 4. Z Berenyi et al, J Electr. Spectr and Rel. Phen. 135, p. 177 (2004) 5. L. Kover et al. Surf. Interf. Anal. 38, p. 569 (2006) 6. S Tougaard, F Yubero Surf Interf Anal. 36, p. 824 (2004) Copyright (c) 2001-2010
			and Francisco Yubero ICMS, Seville, Spain The software is free for non commercial use



1.75E+00

1.40E+00

1.05E+00

6.99E-01

3.49E-01

0.00E+00

0

10



Calculates differential cross sections for 65 depths















#### XPS-Simulation

File Tools Tools exclusively for the L\_K spectrum View Copyright



XPS-Simulation			
File Tools Tools exclusively for the L_K s	pectrum View Copyright		
XPS Simulation     c: [Vista]     C:\     Quases-Tougaard     QUEELS-XPS Simulation     Data	ELF-Si EnergyLossFunction.dll Si2p-Keff-XPS-Electr Si2p-Keff-XPS-hole Si2p-Keff-XPS-Total	Calculate the effective XPS-cross section from Keff(X0) -1. Click Keff(X0) data file in 2. Press Read Keff(X0): Si2p-Keff-XPS-Total Plot and Save individual Keff(X0)	
		3. Press one of the buttons: Main menu	Calculate Total effective XPS cross section for film of thickness End depth End layer = 65
x: .00, y: 0.00E+00			Calculated cross section











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## QUEELS-XPS software package Semi-classical dielectric response model

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ST,FY: SIA 36(2004)824

# **Effect of the static hole – intrinsic excitations**



Yubero, Tougaard: SIA 36(2004)824

the hole contributes little to the surface plasmon but a lot to the bulk plasmon.

#### Two holes left behind... Auger emission



The intensity of the plasmon excited is larger for 2 holes than for 1 hole

Note: Only input in the theory is the dielectric function  $\varepsilon(k,\omega)$ 

 $\varepsilon(k,\omega)$  can be taken

from litterature or

easily be determined from analysis of a REELS spectrum using the QUEELS-ε(k,ω)-REELS software package

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#### Si2p photoemission compared to experiment

Phys. Rev.B56 (1997)1612

![](_page_34_Figure_2.jpeg)

# Variation of emission angle

![](_page_35_Figure_1.jpeg)

Comparison is on an absolute scale

Only input in the theory is  $\varepsilon(k,\omega)$ 

# Variation of emission angle

Yubeo et al Surf Sci 592 (2005) 1

![](_page_36_Figure_2.jpeg)

# Variation of photon energy 225-850 eV

Yubeo et al Surf Sci 592 (2005) 1

![](_page_37_Figure_2.jpeg)

# Variation of photon energy 3000 eV – 5500 eV

![](_page_38_Figure_1.jpeg)

Higher photon energy gives larger intrinsic plasmon (less time to relax and screen the hole potential).

#### Theory

![](_page_39_Figure_3.jpeg)

![](_page_40_Figure_0.jpeg)

Experiment

Cu2p

#### Theory

![](_page_41_Figure_3.jpeg)

Cu 2p - XPS spectra 1.0 -•0° simulations 60° 0.8 75° 82° - 82° 0.6 75° 0.4 **₄**60° **,** F(E) 0.2 0° 0.0 520 540 560 500

![](_page_42_Figure_0.jpeg)

# QUEELS-XPS software package

# **Conclusion:**

The semiclassical dielectric model for photoemission reproduces rather well quantitatively the angular and energy dependence of the energy loss processes

We hope the XPS-model and QUEELS-XPS software package will be useful to gain better quantitative understanding of the photo-excitation process

The software was used for the work presented by N. Pauly this morning

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The software was used in the talk by Nicolas Pauly Tuesday morning. "New parameter to correct for surface and core hole effects in XPS"

![](_page_45_Figure_0.jpeg)